

## ELEMENTARY PARTICLE PHYSICS

FORCES OF NATURE – FUNDAMENTAL INTERACTIONS (PART III) - QCD

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#### Outline:

- A brief **introduction** (history ...)
- The **tools** (accelerators, targets, detectors ... kinematics, ...)
- The **particles** (hadrons, baryons, mesons ...)
- The fundamental particles (quarks, leptons)
- The forces (gravitation, nuclear, weak forces)
- The *fundamental* interactions (strong and electro-weak IA)
- The Standard Model of EPP
- Physics **Beyond the Standard Model** (BSM)
- Spin-offs Applications of EPP



#### Prelude

#### **History** – the genesis of QCD (I)

In 1971, **M. Gell-Mann** and **H. Fritzsch** found a solution of the statistics problem (remember the  $\Omega^{-}$ ): they considered **nine quarks** (as others had done before) (<u>note</u>: **only u,d,s** known at that time), but assumed that the three quarks of the same type had a **new conserved quantum number**, called "color":





#### Prelude

#### History – the genesis of QCD (II)

It turns out that **color** is the **charge** of the **strong interaction** (color charge – like the electric charge for the electromagnetic interaction):



Quarks, which carry color (anti-quarks – anti-color), interact via the exchange of "**gluons**" (analogous to photons which mediate the electro-magnetic interaction).



#### Prelude

#### **History** – the genesis of QCD (III)

"Quantum chromodynamics" (QCD) is the theory of the strong interaction between quarks and gluons; it has been modeled in analogy to QED; the dynamics of the quarks and gluons are controlled by the QCD *Lagrangian* and visualized by Feynman diagrams; one basic diagram is:



QCD is an important part of the Standard Model of particle physics



#### **Basics**

**Facts** – the QCD exchange bosons – colored gluons (I)

Gluons act as the exchange particle of the strong force between **colored quarks**; gluons are carrying **both color and anti-color**:



There are **nine possible combinations** of color and anti-color in gluons: red anti-red, red anti-blue, red anti-green, blue anti-red, blue anti-blue, blue anti-green, green anti-red, green anti-blue, green anti-green.



#### **Basics**

**Facts** – the QCD exchange bosons – colored gluons (II)

Analogous to the  $SU(3)_F$  flavor (for the 3 light quarks (u,d,s) – see: hadron multiplets) the 3 colors can form an  $SU(3)_C$  symmetry group with a **color octet** and a **singlet**:



with the singlet state not realized in Nature



**Facts** – the QCD exchange bosons – colored gluons (III)

Since **gluons** carry color – anti-color they not only **interact** with colored quarks, but also **among themselves** (in contrast to photons, which do not carry electric charge):



The "**gluon self interaction**" (gluons themselves take part in strong IA) is the reason for **fundamental differences** between **QED** and **QCD** 



#### **Basics**

- **Facts** the QCD coupling constant (I)
  - The strong coupling constant  $\alpha_{s}$  determines the strength of the strong interaction
  - In **QED** the bare electric charge is screened by a cloud of virtual (e<sup>+</sup>e<sup>-</sup>) pairs, leading to the "running"  $\alpha$  (increase with Q), in **QCD** there are two such effects:
  - screening of the color charge by virtual quark-antiquark pairs
  - anti-screening by a cloud of virtual gluons





#### **Basics**

#### **Facts** – the QCD coupling constant (II)

The two effects are leading to the **"running" of**  $\alpha_s$ , and it turns out: **anti-screening dominates** – the effective **color charge increases with distance**, and **decreases with energy** (momentum transfer):





#### **Basics**

#### **Facts** – the QCD coupling constant (III)

The "running" of  $\alpha_s$  is very well **established experimentally**:



→ 2 regimes: perturbative ( $\alpha_s << 1$ ), non-perturbative, strong ( $\alpha_s \sim 1$ ) QCD



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#### **Facts** – the QCD coupling constant (IV)

In the high-energy/small-coupling regime, the quarks are essentially behaving as free particles: this is called "**asymptotic freedom**"; for lowenergy, the coupling becomes so large that quarks cannot be separated: (quark) "**confinement**":



Asymptotic freedom of QCD was discovered by **D. Gross**, **F. Wilczek** and **D. Politzer** (1973)  $\rightarrow$  NP 2004. Quark confinement of QCD has yet to be proven from first principles



#### **Basics**

**Facts** – the QCD exchange bosons – reality of gluons (I)

The process of "gluon bremsstrahlung":

Example: e<sup>+</sup>e<sup>-</sup> annihilation to a photon; generation of a qq-pair fragmenting into hadrons and appears in the detector as two back-toback hadron jets. Outgoing quarks can also radiate a gluon, creating a third hadron jet in the same plane as the other two



Discovered at DESY (Hamburg, Germany) in 1979



#### **Basics**

Facts – the QCD exchange bosons – reality of gluons (II)

The process of "gluon bremsstrahlung":

Example: 2- and 3-jet events





#### **Basics**

#### Facts - the QCD exchange bosons - reality of gluons (III)

The angular distribution of the jets depends on the gluon spin:



Experiments find that **gluons** are **vector bosons** (spin-1) as expected for exchange particles



#### **Basics**

Facts – the basic vertices of QCD

Gluon absorption and emission:



#### Quark – anti-quark annihilation and pair production:

 $\overline{q}$  g g g  $\overline{q}$ 

<u>Note</u>: the vertices do not yet represent physical processes (require a combination)



#### **Basics**

#### Facts – Feynman's partons

R. Feynman postulated (in 1969) that protons ( $\rightarrow$  hadrons) were made of pointlike constituents, he called "**partons**". Later, it was recognized that partons describe the same objects now more commonly referred to as **quarks** and **gluons**:



<u>Note</u>: When probed at smaller scales, e.g. in **DIS** (see below), protons seem to contain more and more **partons** (quarks and gluons), represented here as colored blobs.



#### **Examples**

#### Hadronic processes – generic reaction

A virtual photon ( $\gamma^*$ ) (produced, e.g., in e<sup>+</sup>e<sup>-</sup> annihilation) produces a **quark – anti-quark pair**:



Not only **electric charge** but also **color charge** of quarks and **gluons** are involved



#### Examples

#### Hadronic processes – Drell-Yan process (I)

In **pp-collisions** the so called "**Drell-Yan**" (DY) process occurs: a quark from one proton and an antiquark from the other proton (note: a "sea quark"!) annihilate into a **virtual photon** ( $\gamma^*$ ). The photon can split into a lepton and its antiparticle partner, for example into an e<sup>+</sup>e<sup>-</sup> or  $\mu^+\mu^-$  pair, provided the  $\gamma^*$  energy is sufficient:



(DY can also happen via an intermediate Z-boson (weak IA, later ...))



#### Examples

#### Hadronic processes – Drell-Yan process (II)

The Drell-Yan process is not so simple due to the **complexity of the proton**: real collisions also include the **remnants** of the scattered

protons:



<u>Notes</u>: (i) complex final state; identification of  $(\mu\mu)$ -pair simple (ii) many more graphs (loops ...)



#### **Examples**

#### Hadronic processes – Drell-Yan process (III)

<u>Example</u>:  $M_{u+u}$  spectrum in pp-collisions at LHC



The rightmost peak at about 90 GeV (~90 times the proton mass!) is a peak corresponding to the production **Z bosons** (→ weak IA). The other peaks represent the production of well-known particles that have decayed into a muon-antimuon pair.



#### **Examples**

#### Hadronic processes – Drell-Yan process (IV)

The use of **proton** – **anti-proton** collisions is advantageous, because the anti-proton contains valence anti-quarks:



→ In addition: polarization as an additional/new degree of freedom (<u>our project</u>: production of **polarized anti-protons** ...)



#### **Examples**

#### Hadronic processes – Drell-Yan process (V)

DY in **proton** – **anti-proton** collisions may also be used to search and investigate **exotics** (like tetraquarks):



→ Possible science case for **PANDA at HESR** (FAIR) ...



#### **Examples**

#### Hadronic processes – deep-inelastic scattering (DIS) (I)

Make things simpler: use an electromagnetic probe (i.e **high-energy scattering** of **charged leptons** on nucleons (avoid complexity of second nucleon):





### FUNDAMENTAL INTERACTIONS – QCD Examples

#### Hadronic processes – deep-inelastic scattering (DIS) (II)

Elastic electron scattering (eN  $\rightarrow$  eN; form factors) and excitation of nucleon resonances (eN  $\rightarrow$  eN\*  $\rightarrow$  eN X) already discussed; at **higher energy** (smaller virtual photon wave length), on probes the **internal constituents**; at same time strong coupling constant becomes small  $\rightarrow$ use of **perturbative approximation** (like in QED):





#### **Examples**

Hadronic processes – deep-inelastic scattering (DIS) (III)

DIS is **elastic scattering on a quark** inside a nucleon:



Since quarks are supposed to be point-like, the corresponding formfactors should be constant ( $\rightarrow$  "**Bjorken scaling**", J. Bjorken, 1968; Feynman's point-like partons)  $\rightarrow$  this inspired QDC



#### **Examples**

#### Hadronic processes – deep-inelastic scattering (DIS) (IV)

Bjorken scaling (independence of "**structure functions**" on momentum transfer) is experimentally observed:

Example: HERA (DESY) e<sup>+</sup>p and e<sup>-</sup> p data





#### Examples

#### Hadronic processes – deep-inelastic scattering (DIS) (V)

- Bjorken scaling is not exact; deviations from strict scaling is required in quantum field theory; **QCD can predict** the detailed form of **violations of the scaling behavior** of the relevant physical quantities.
- Quarks inside a nucleon have a **momentum distribution** each one carries a varying fraction of the energy/momentum of the nucleon; the momentum distribution can be determined by looking at the scattered

electrons:



Expt'l finding: quarks carry about 50% of the proton momentum; gluons carry another ~50%



#### Examples

Hadronic processes – deeply-virtual Compton scattering (DVCS)

Even simpler: in "**Deeply Virtual Compton Scattering**" (DVCS), a highenergy electron probes a nucleon by exchanging a **virtual photon** with the quarks inside. The final-state **real photon** carries information about the nucleon structure:



Experimental problem: separation of "Bethe-Heitler" background



#### **Examples**

#### Hadronic processes – nucleon spin (I)

In the simple CQM, the spin  $S_N = \frac{1}{2}$  of the nucleon is just the vector sum of the 3 quark spins of  $S_q = \frac{1}{2}$  (with two parallel and the 3<sup>rd</sup> one anti-parallel):



<u>But</u>: the nucleon is a much more complex object (valence and sea-quarks, gluons)  $\rightarrow$  what is their contribution?



#### Examples

#### Hadronic processes – nucleon spin (II)

Experimentally, it is found that the **quark spin contributes about 30%** to the spin of the nucleon ( $\rightarrow$  "**nucleon spin crisis**"); a major topic of expt'l particle physics is to find the missing part, believed to be carried either by **gluon spin**, or by **gluon** and **quark orbital angular momentum**:





#### **Examples**

#### Hadronic processes – from the qq interaction to the nuclear force

The strong interaction binds quarks inside nucleons, and the **residual strong force** (nuclear force) binds nucleons in nuclei:





#### Examples

Hadronic processes – production of the Higgs boson

The dominant **Higgs boson production mechanism** (~88%) at the elementary level by **gluon fusion** and a quantum loop process involving a **top quark**; discovery at the **LHC** in pp collisions (2012):



Talk about (wish for) a **Higgs-factory** ... highest priority for the particle physics community: CLIC, FCC, ILC, ...



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#### Outlook

Hadronic processes - new (upcoming) experimental facilities

**EIC** (Electron-Ion Collider) – Brookhaven (USA)



FAIR (Facility for Antiproton and Ion Research) – GSI (Germany)





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#### Summary

#### Quantum Chromodynamics (QCD) – main features

#### Confinement

- At large distances the effective coupling between quarks is large, resulting in confinement.
- Free quarks are not observed in nature.
- Asymptotic freedom
  - At short distances the effective coupling between quarks decreases logarithmically.
  - Under such conditions quarks and gluons appear to be quasi-free.
- (Hidden) chiral symmetry
  - Connected with the quark masses
  - When confined quarks have a large dynamical mass constituent mass
  - In the small coupling limit (some) quarks have small mass current mass





### **THE FORCES**

#### That's it for today





The coupling between quarks and gluons depends strongly on the energy scale of the process. The same is true for the masses of the quarks. This effect – the so-called "running" of the strong coupling constant and the quark masses – is described byquantum chromodynamics (QCD). The experimental verification is both an important test of the validity of QCD and an indirect search for unknown physics





